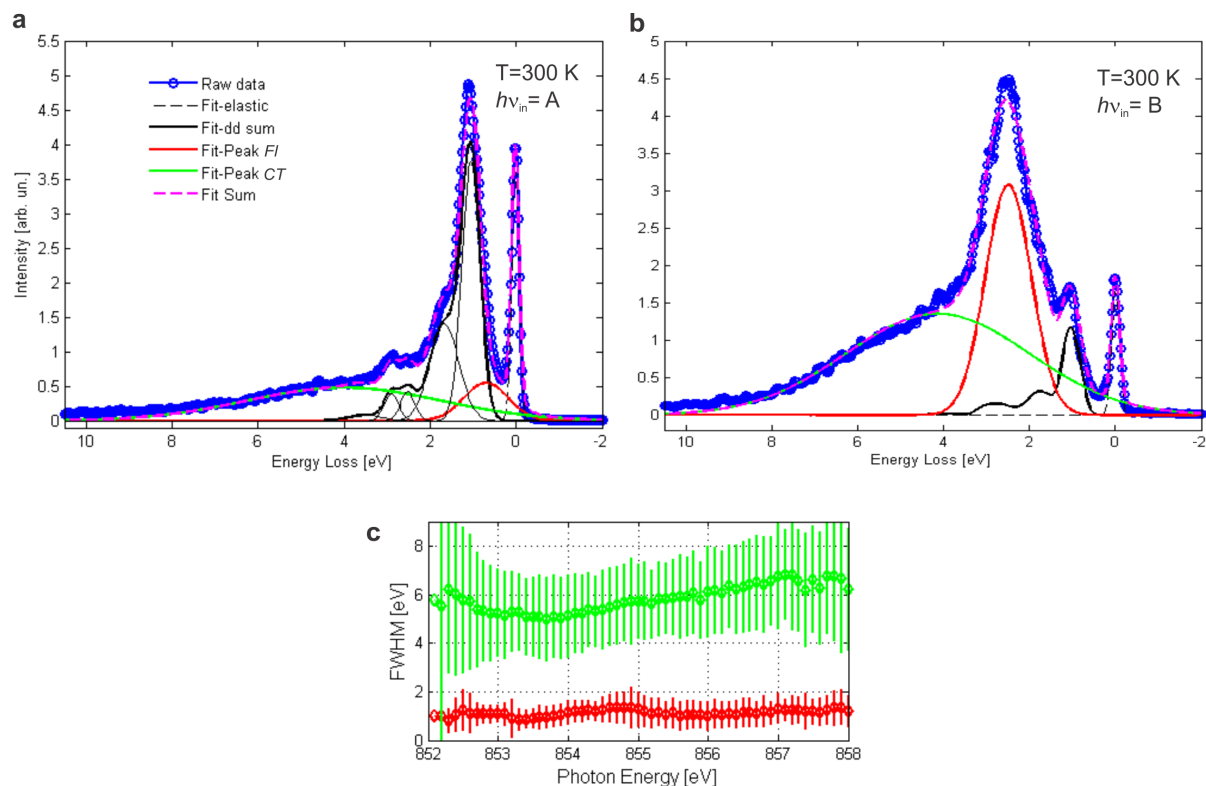
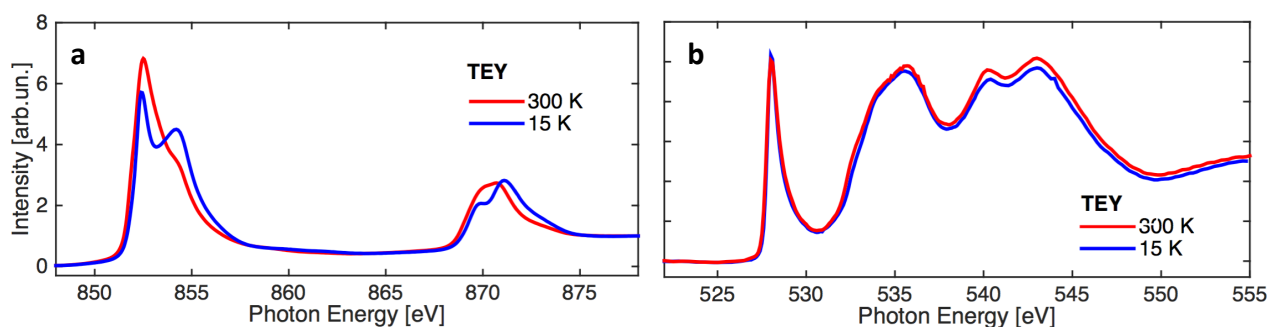




Supplementary Figure 1 | Details of fitting analysis for RIXS data at 300 K. **a**, RIXS spectrum measured at $h\nu_{in}=A$. The thin solid lines refer to the Gaussians used to fit the individual contributions of the spectrum: the elastic line (black dashed line), the intra-band dd -excitations (thin black solid lines), a delocalized contribution with a FWHM of 1eV (Peak FI in red) and a charge-transfer contribution with a FWHM of 6 eV (Peak CT in green). The black thick solid line represents the sum of the dd -contributions. The magenta dashed line represents the sum of all the Gaussian contributions. **b**, RIXS spectrum and similar decomposition as in (a) for $h\nu_{in}=B$. **c**, FWHM of FI (red) and CT (green) peaks resulting from the fitting analysis, as a function of the Photon Energy at 300 K. The error bars of the FWHM parameters are evaluated using the least square fitting routine and expressed in (s.e.d.).



Supplementary Figure 2 | X-ray Absorption. **a**, Ni $L_{3,2}$ X-ray Absorption of $NdNiO_3$ and, **b**, O K XAS in Total Electron Yield. XAS spectrum at 300K (15K) shown in red (blue).

Parameter	Value	Description
$F_{dd}^{(2)} / F_{dd}^{(4)}$	9.787 / 6.078	Intra-atomic Slater (Coulomb) integrals of Ni 3d electrons
$F_{2p3d}^{(2)}$	7.392	Intra-atomic Slater (Coulomb) integrals of Ni 2p and 3d electrons
$G_{2p3d}^{(1)} / G_{2p3d}^{(3)}$	5.498 / 3.126	Intra-atomic Slater (exchange) integrals of Ni 2p and 3d electrons
ζ_{2p} / ζ_{3d}	11.507 / 0.083	Spin orbit coupling of the indicated shell
$10Dq$	0.65	Ionic crystal field splitting
V_{eg}	2.06	Hopping integral for e_g symmetry
$V_{t_{2g}}$	1.21	Hopping integral for t_{2g} symmetry
Δ	4.20	Charge transfer energy (to center of band)
U_{dd}	6.50	Monopole part of on-site Coulomb repulsion
U_{pd}	7.50	Monopole part of 2p core hole potential
H_{ex}	0.10	Local exchange field
W	5.40	Width of ligand valence band
N	25 points	Number of discretization points for ligand band

Supplementary Table 1 | Parameters used for the impurity model calculations. Values are in units of eV, unless otherwise indicated.

Supplementary Note 1:

RIXS Spectra fitting analysis

The rich structures of the RIXS spectral lineshapes, represented in Figure 3 (main text) at 15 K and in Supplementary Figure 1 at 300 K (for complementarity), have been analysed and fit by means of spectral decomposition based on Gaussian curves [1]. We use for this analysis the full dataset of the RIXS energy maps. The main prominent RIXS contribution at the XAS resonance $\hbar\nu_{in}=A$ (852.4 eV) stems from the localised dd -excitations in the range between 1-3 eV. Examining all the RIXS spectra corresponding to incoming photon energies between 852 eV to 859 eV (across the full Ni L_3 -edge), we notice that the dd -excitations do not significantly change in their total profile, but mainly in their intensity. Therefore we determine the dd -profile by referring to the resonance spectrum at 852.4 eV [see Suppl. Figure 1 (a)]. Four Gaussian curves (thin black lines) have been tuned to match the sharp peaks emerging from the spectrum in the 1-3 eV energy loss range. Their sum leads to the total dd -profile (thick black line). At the same time, the broad residual spectral weights remaining at around 0.7 eV and 4 eV have been fit

with two other Gaussian curves (red and green lines), shallower in intensity but much broader than the previous ones.

Taking advantage of the constant *dd*-profile throughout the Ni L₃-edge, we proceed to decompose all the RIXS spectra between 852 eV to 859 eV as the sum of three Gaussians - one for the elastic line (black dashed line), one for *FI*-peak (red solid line) and one for *CT*- peak (green solid line) – and a properly rescaled *dd*-profile (black thick line) [refer to Suppl. Figure 1 (a-b) for 300K and to Figure 3 (a) from main text for 15K]. The decomposition is done by using a fitting analysis based on the least squares method, and the components listed above define the number of free parameters required for this fitting: (1) the intensity rescaling factor for the *dd*-profile (normalized to unit area); (2-7) the intensity, the peak position and the FWHM of the two Gaussians used respectively to track the *CT* and *FI* peaks; (8) the intensity of the Gaussian used to fit the elastic peak (the position is by definition at 0eV and its width mostly determined by the experimental resolution).

The extracted fitting parameters as a function of the incident photon energy, and their respective errors, are presented in Figure 3 (main text) and therein discussed in terms of intensity [Figure 3 (b-c)] and peak position [Figure 3(d-e)]. In addition, we plot in Suppl. Figure 1 (c) the extracted FWHM for the *CT* and the *FI* peaks. We note that these values remain reasonably constant ($\pm 10\%$) in the full energy range and average at 1 eV for *FI*- and 6 eV for the *CT*-peak. Such a result further supports the validity of the presented decomposition analysis, allowing characterizing the energy dispersion of the *CT* and *FI* contributions as a function of the incident photon energy.

Supplementary Note 2:

X-ray Absorption in Total Electron Yield

X-ray absorption measurements in Total Electron Yield mode (TEY) have been performed for NdNiO₃. Spectra acquired for 300K and 15K are presented in Suppl. Figure 2, both at Ni L and O K edges. Given the strong surface-sensitivity of the TEY signal, we used in the main part of the

paper the Partial Fluorescence Yield (PFY) obtained by integrating the RIXS spectra at the Ni L₃-edge, to ensure the bulk-sensitivity. The XAS in PFY mode displayed in Figure 2(a) of the main text is in good agreement with previous data published in the literature [2, 3]. This confirms the good quality of our sample. Concerning the O K XAS spectra, our measurements displayed in Supplementary Figure 2(b) are also in good agreement with the main features observed in the literature, on the bulk NdNiO₃ sample [2, 3], thus showing: 1) a very strong pre-peak around 528 eV which is proportional to the covalent degree in the system; 2) a small spectral weight increase at the prepeak and a decrease at the 536 eV peak, while going to the insulating phase [3].

Supplementary Note 3:

Cluster and Crystal Filed calculations for orbital excitations in NdNiO₃

The spectral shape of the *dd* excitations is strongly determined by the local Ni 3*d* orbital energies and occupations, and accordingly provide direct insight into the Ni valence. To analyze the observed *dd* excitations, we first computed the RIXS response of a typical 3*d*⁷ cluster model, which has been used to interpret the XAS of the nickelates in the past [3]. In this model, a 3*d*⁷ Ni impurity hybridizes with a fully occupied oxygen octahedron, leading to a mixed, many-body ground state wavefunction of the form $|\psi_0\rangle = \alpha|d^7\bar{L}^0\rangle + \beta|d^8\bar{L}^1\rangle + \gamma|d^9\bar{L}^2\rangle$, where $\bar{L}$ denotes a ligand hole. We took all model parameters (multipole Coulomb interactions, crystal field splitting, hopping integrals, charge-transfer energies) from Ref. [3]. RIXS spectra were calculated using the Kramers-Heisenberg equation. The spectra were broadened with Lorentzian and Gaussian lineshapes to account for lifetime and instrumental broadening effects, respectively. We find that the computed RIXS spectra show *dd* excitations which strongly disagree with experiment, including a low energy low-spin to high-spin excitation near 0.25 eV energy loss which is not found in experiment, and other multiplets which deviate strongly from those observed experimentally.

Next, we calculated the expected dd excitations from a Ni $3d^8$ configuration having the local O_h symmetry (approximately) present in the nickelates. We find that even using a simple crystal field model, excellent agreement with the experimental dd excitations is obtained, by employing model parameters (multipole Coulomb interactions, crystal field splitting) which have been used in studies of NiO [4]. In particular, we have for the Slater integrals $F^2 = 8.063$ eV and $F^4 = 5.699$ eV, and for the crystal field splitting a value of $10D_q = 1.05$ eV. Including hybridization effects would add charge-transfer excitations, but would not strongly alter the dd excitations after a renormalization of the crystal field energy [4].

Supplementary Note 4:

Singe Impurity Anderson Model calculations

The Single Impurity Anderson Model (SIAM) calculations used (in Figure 6 and related discussion in the main manuscript) to analyze the behaviour of local dd and nonlocal charge-transfer excitations include full Ni $3d$ -shell multiplets and coupling to a fully occupied, ligand O $2p$ -like valence band. The model parameters used are shown in Supplementary Table 1 and are typical for a NiO-type system, agreeing well with both *ab-initio* NiO values [5] and those fitted to experiment [6,7]. Note, however, that the precise values for the parameters are not important for the analysis presented in this work. For the $3d^8$ impurity configuration considered here, the basis sizes are relatively small (of order 10^3), thus easily allowing full diagonalization of the Hamiltonian matrices for each irreducible representation. XAS and RIXS spectra were then calculated using Fermi's Golden Rule #2 and the Kramers-Heisenberg equation, respectively [8]. The spectra were broadened with Lorentzian and Gaussian lineshapes to account for lifetime and instrumental broadening effects, respectively.

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