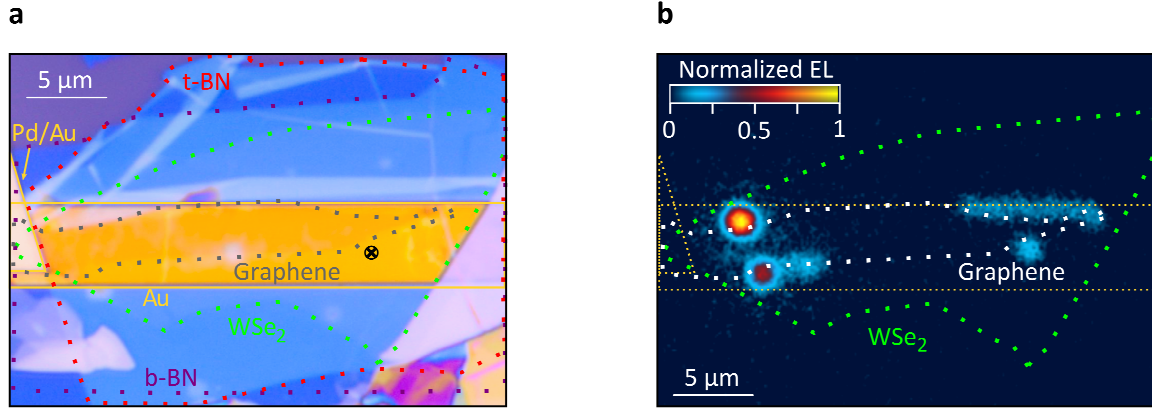


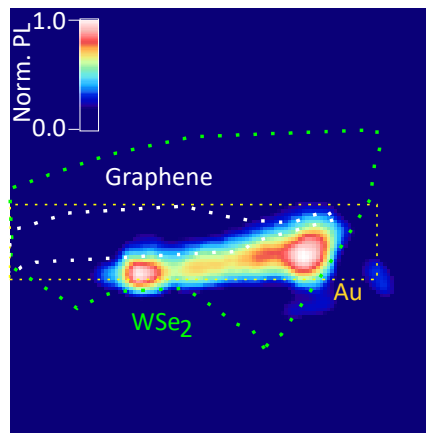
Supplementary Information

Electroluminescence from multi-particle exciton complexes in transition metal dichalcogenide semiconductors

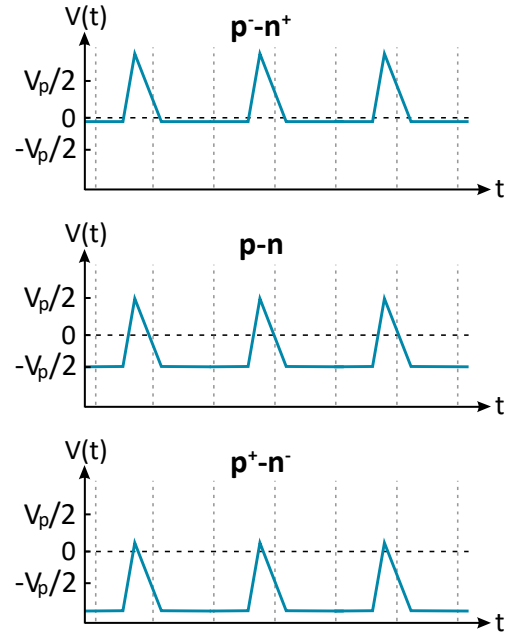
Paur et al.



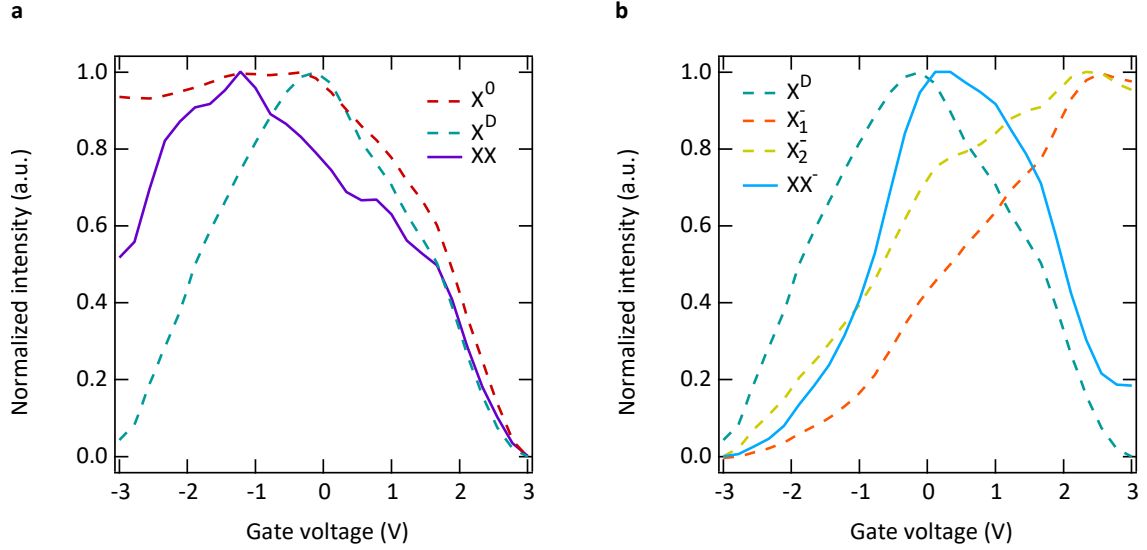
Supplementary Figure 1 | Electroluminescence image of WSe₂ sample at 300 K. **a**, Optical microscopy image of the sample. **b**, False-color EL image of the sample. The EL emission is generated near the interface between graphene and monolayer WSe₂. However, the light emission is not homogeneously generated along this interface, because part of the graphene/WSe₂ junction is not gated and possibly also because of an inhomogeneous contact between graphene and WSe₂. Also note that in this particular device the WSe₂ flake bends over the bottom electrode edge, facilitating the detection of the emission from dark excitons with out-of-plane dipole (see Figs. 2c and 2e in the main text).



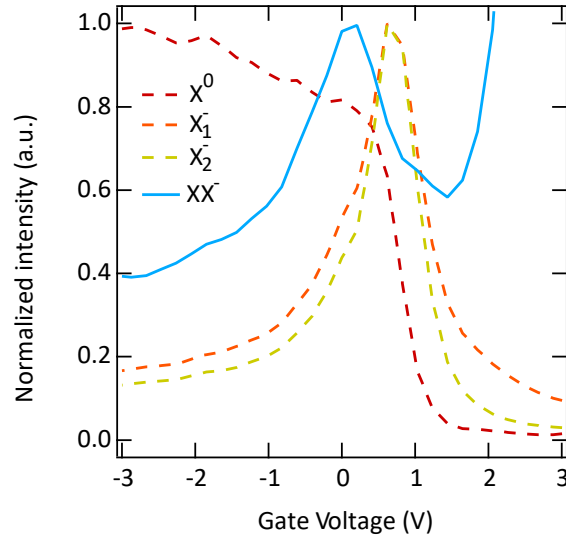
Supplementary Figure 2 | Photoluminescence map at 300 K (WSe₂ sample). Photoluminescence from monolayer WSe₂ at room temperature (excitation $\lambda = 532$ nm, $P_d = 8 \times 10^3$ W cm⁻²). The PL is quenched on graphene.



Supplementary Figure 3 | Schematic pulse diagrams. Top: electron-rich (p^-/n^+); Middle: balanced (p/n); Bottom: hole-rich (p^+/n^-) carrier injection.

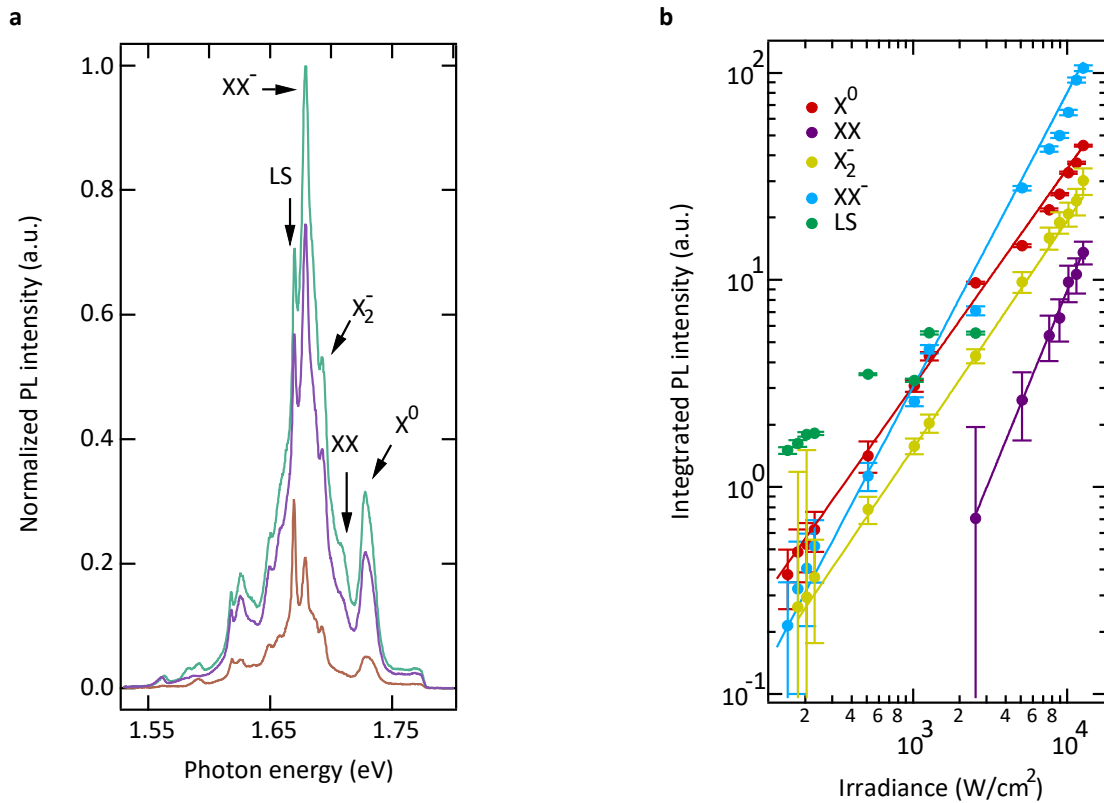


Supplementary Figure 4 | PL emission from exciton complexes as a function of external electric field in WSe₂. Line-cut of the normalized PL intensity of different exciton species. **a**, Biexcitons (purple) **b**, charged biexcitons (blue) as a function of doping level. The behavior of the different peaks with the back voltage is in good agreement with recently reported results¹⁻⁴.



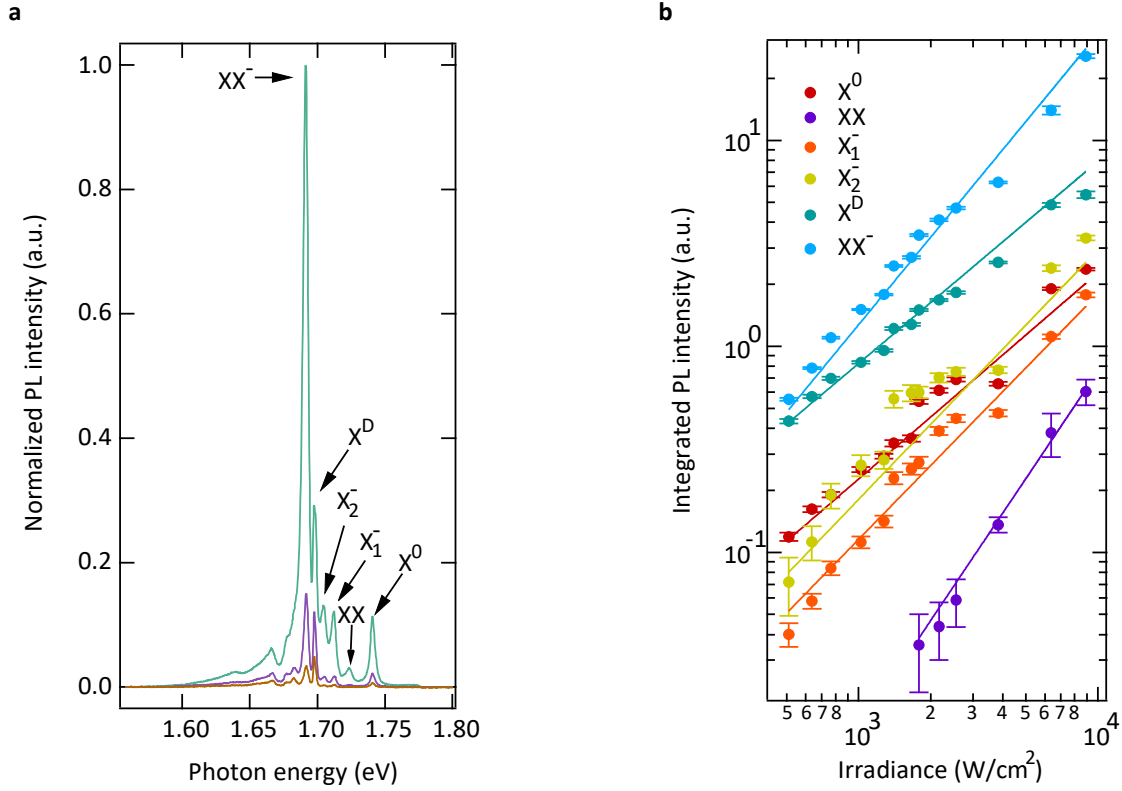
Supplementary Figure 5 | PL emission from exciton complexes as a function of external electric field in WS₂.

Line-cut of the normalized PL intensity of different exciton species as a function of doping level. In the highly n-doped regime ($V_G > 2$ V), where X^0 (red), X_1^- (orange), X_2^- (yellow), XX^- (blue) species vanish, the next charging state of the trion may be observed, as in WSe₂³.

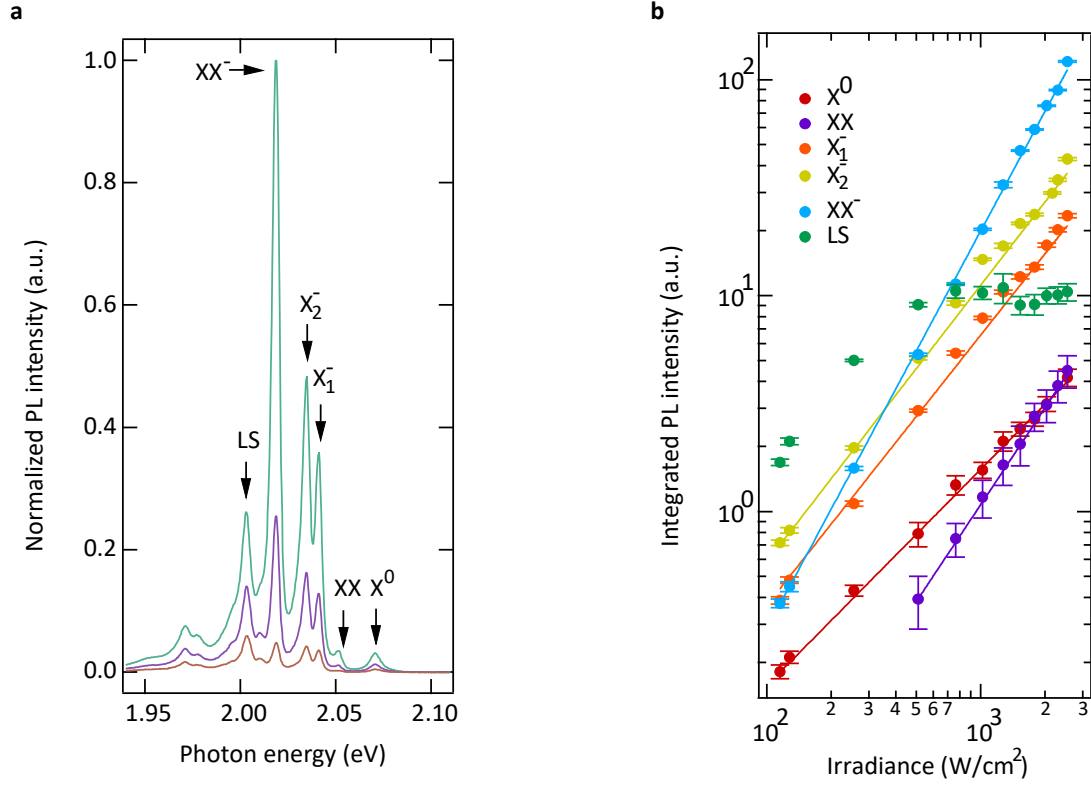


Supplementary Figure 6 | Power-dependence of the PL intensity in WSe₂. **a**, PL spectra of a WSe₂ monolayer for different excitation intensities ($P_d = 2.5 \times 10^3$ Wcm⁻² to 1.2×10^4 Wcm⁻²). **b**, Logarithmic plot of the integrated PL intensity as a function of the excitation intensity. Circles represent the experimental data, while lines

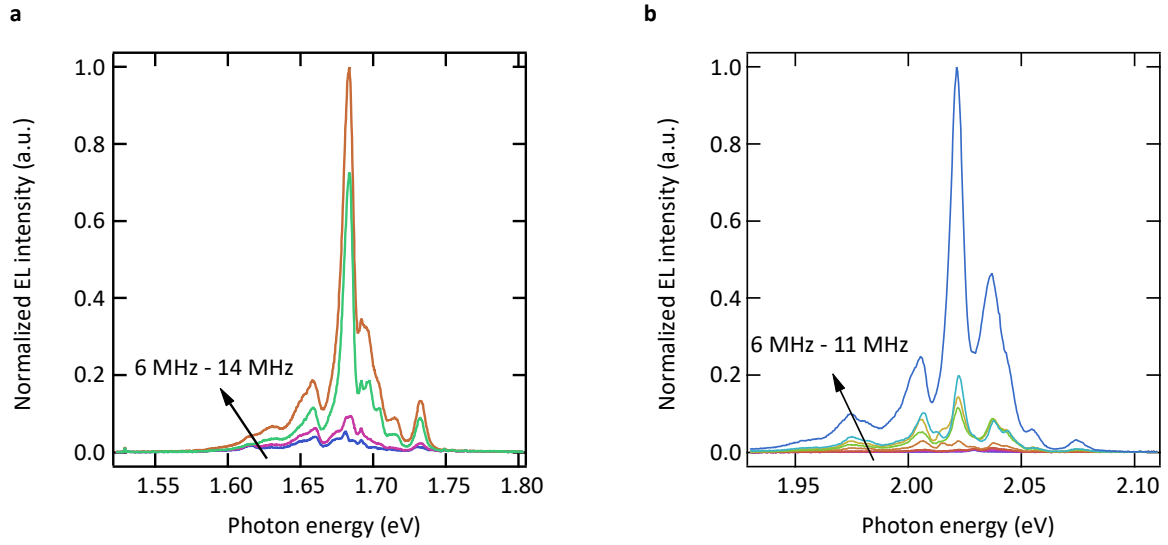
represent a power law fit, $\propto P^\alpha$, with exponents of X^0 ($\alpha = 1.01$), XX ($\alpha = 1.82$), X_2^- ($\alpha = 1.10$), XX^- ($\alpha = 1.44$). The peak LS ($\alpha = 0.6$), centered at an energy of 1.6696 eV, is most probably due to a localized state, since the dependence on excitation intensity is sublinear^{5–10}. The power-dependence measurements are performed without gate voltage (intrinsic regime) at the position marked with a cross in Supplementary Figure 1a.



Supplementary Figure 7 | Power-dependence of the PL intensity in another WSe₂ sample. **a**, PL spectra of a WSe₂ monolayer for different excitation intensities ($P_d = 5.0 \times 10^2$ Wcm⁻² to 6.3×10^3 Wcm⁻²). **b**, Logarithmic plot of the integrated PL intensity as a function of the excitation intensity. Circles represent the experimental data, while lines represent a power law fit, $\propto P^\alpha$, with exponents of X^0 ($\alpha = 1.01$), XX ($\alpha = 1.73$), X_1^- ($\alpha = 1.19$), X_2^- ($\alpha = 1.21$), X^0 ($\alpha = 0.99$), XX^- ($\alpha = 1.41$). The power-dependence measurements are performed without gate voltage (intrinsic regime).



Supplementary Figure 8 | Power-dependence of the PL intensity in WS₂. **a**, PL spectra of a WS₂ monolayer for different excitation intensities ($P_d = 5.1 \times 10^2 \text{ Wcm}^{-2}$ to $2.5 \times 10^3 \text{ Wcm}^{-2}$). **b**, Logarithmic plot of the integrated PL intensity as a function of the excitation intensity. Circles represent the experimental data, while lines represent a power law fit, $\propto P^\alpha$, with exponents of X^0 ($\alpha = 1.01$), XX ($\alpha = 1.49$), X_1^- ($\alpha = 1.25$), X_2^- ($\alpha = 1.28$), XX^- ($\alpha = 1.84$). The peak LS ($\alpha = 0.2$) at an energy of 2.006 eV, is most probably due to a localized state or phonon-assisted emission^{11,12}. The power-dependence measurements are performed without gate voltage (intrinsic regime) on the position marked with a cross in Figure 1b of the main text.



Supplementary Figure 9 | EL intensity as a function of the pulse frequency. **a**, WSe₂ EL spectra at frequencies ranging from 6 MHz to 14 MHz. **b**, WS₂ EL spectra at frequencies ranging from 6 MHz to 11 MHz. The emission intensity increases with frequency, in agreement with similar measurements performed at room temperature¹³. Furthermore, all the emission peaks broaden at higher frequencies.

	X^0	XX	X_1^-	X_2^-	X^D	XX^-	L_1
WSe₂							
EL	1	1.44	-	-	1.03	1.75	-
PL	1.01	1.82	1.19*	1.1	1.08	1.44	0.6
WS₂							
EL	1	2.33	1.32	1.33	-	2.31	-
PL	1.00	1.49	1.25	1.28	-	1.84	0.2

Supplementary Table 1 | α coefficient values obtained from the power-law fits in EL and PL for the different exciton species. *This value is obtained from measurements performed in a WSe₂ sample different from that presented in the main text and shown in Supplementary Figure 7.

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