

Supplementary Information for
“Direction-dependent Dynamics of Colloidal
Particle Pairs and the Stokes-Einstein
Relation in Quasi-Two-Dimensional Fluids”

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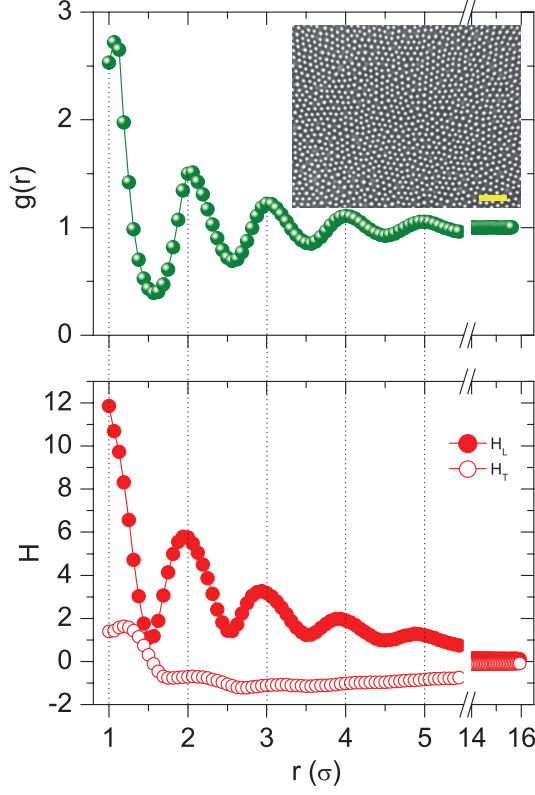


Fig. S1: Connection between hydrodynamics and underlying structure of the fluid. **a**, The sample pair correlation function, $g(r)$, with r expressed in units of particle diameter (σ), at $\phi = 0.61$. Inset: one quarter of a typical field-of-view of the colloidal fluid at $\phi = 0.61$ with scale bar of $5 \mu\text{m}$. **b**, H_L (solid red), H_T (open red) evaluated for $t = 0.5 \text{ s}$ versus r for quasi-two-dimensional colloidal fluids at $\phi = 0.61$. The vertical lines show that the peaks of $g(r)$ and $H_L(r)$ coincide.

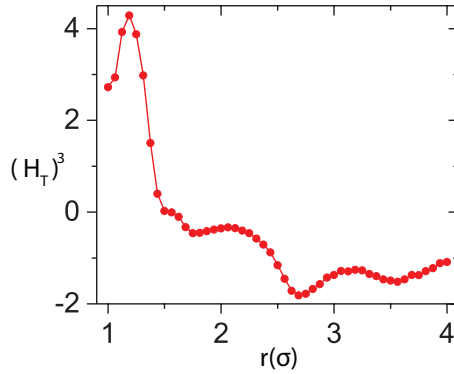


Fig. S2: Revealing spatial positions of maxima and minima in H_T . $(H_T)^3$ versus $r(\sigma)$ at $\phi = 0.61$ confirming spatial phase lag of 0.25σ between H_L and H_T . While H_L peaks at $r = \{1, 2, 3, \dots\}\sigma$, H_T peaks at $r = \{1.25, 2.25, 3.25, \dots\}\sigma$.

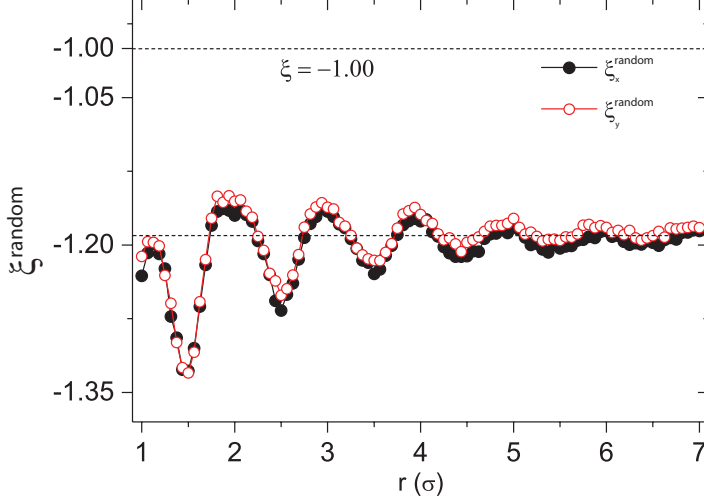


Fig. S3: Stokes-Einstein exponent measured for displacements in random (lab frame) directions. We measured displacement particles in particle-pairs along two randomly chosen orthogonal directions, x and y , in the lab frame. Then we followed the same prescription as described in the main text for the L and T directions to measure the Stokes-Einstein exponent, ξ , along the x and y directions. Notice that the resulting exponents, along randomly chosen directions, vary spatially with r but are *not anisotropic*.

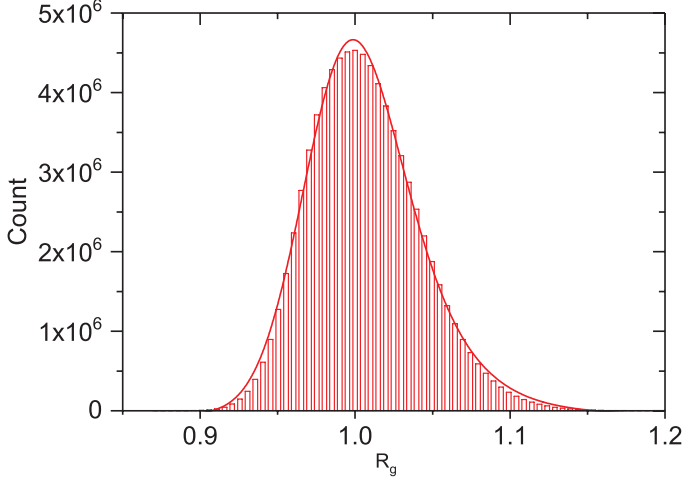


Fig. S4: Ensuring quasi-2D nature of the experimental cell. Histogram of the apparent diameter of the particles, σ^{apparent} , measured from spatial intensity distribution of the tracked particles in the field-of-view from all the images at $\phi = 0.61$. σ^{apparent} is normalized by its most-probable value. The solid line is Gaussian fit to the distribution. The polydispersity in σ^{apparent} , and hence, the size of the particles measured from the distribution $\sim 3\%$, which is comparable to the size polydispersity quoted by the manufacturer ($< 5\%$). This confirms that particles are indeed confined to a quasi-2D plane over the experimental field-of-view. Note, if the height of the cell was larger than the particle diameter, a significant variation in the intensity distribution of particles, and hence, their σ^{apparent} due to the freedom to move out of the focal plane would have been observed. Note, the measured polydispersity could be due to either size polydispersity or height variation or a combination thereof.

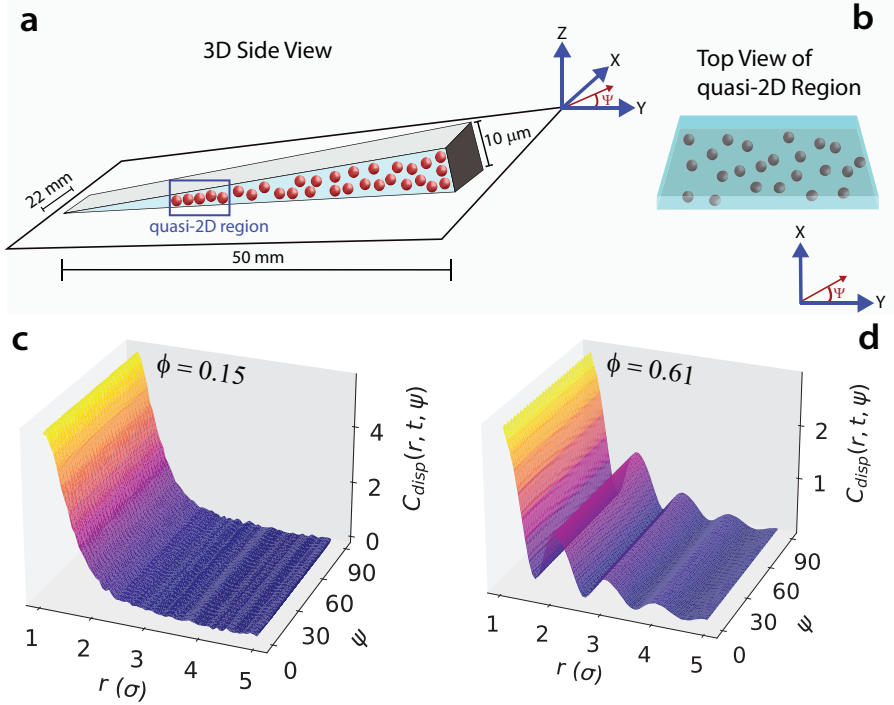


Fig. S5: Schematic of the experimental cell and displacement correlation between colloids in colloid-pair in the lab frame. Schematic of **a**, 3D side and **b**, top quasi-2D region view of the experimental wedge-shaped cell with the wedge along the y -axis in the lab frame. To ensure that choice of the experimental cell does not influence the dynamics measurement of colloids in colloid-pair, we have computed displacement correlation function $C_{disp}(r, t)$ at various angle ψ with respect to the wedge-direction (y -axis in the lab frame). We define $C_{disp}(r, t, \psi) = \langle \Delta \mathbf{r}_i^\psi(\mathbf{r}', t) \cdot \Delta \mathbf{r}_j^\psi(\mathbf{r}' + \mathbf{r}, t) \rangle$. Here, $\Delta \mathbf{r}_i^\psi$ are the displacements of i^{th} particle resolved in a direction making an angle ψ with the y -direction in the lab frame. The three-axes plot of $C_{disp}(r, t, \psi)$ at **c**, $\phi = 0.15$ and **d**, $\phi = 0.61$. Absence of any variation in $C_{disp}(r, t)$ with ψ suggests dynamics revealed in the main text for colloids in colloid-pairs body-frame are due to unique features of hydrodynamic interactions in quasi-2D confinement and rules out its dependence on the choice of the shape of the experimental cell.

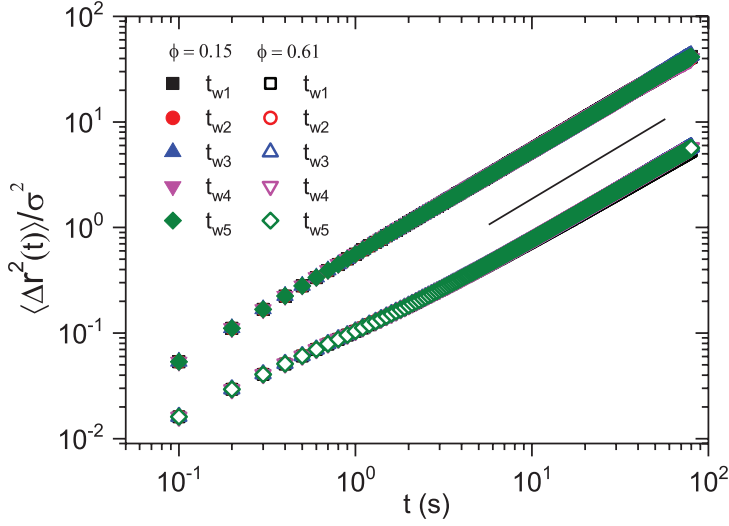


Fig. S6: Ensuring equilibration of samples. To check that colloidal fluids at all ϕ s have equilibrated, the video microscopy data at each ϕ for entire experimental duration was split into 5 equal non-overlapping time windows, t_{w1} to t_{w5} . The plot shows mean-squared displacement $\Delta r^2(t)$ versus t for each of these time windows at $\phi = 0.15$ (solid symbols) and at $\phi = 0.61$ (open symbols). Clearly, the dynamics is not dependent on the time-windows chosen and hence the colloidal fluid at each particle packing area fraction has equilibrated.

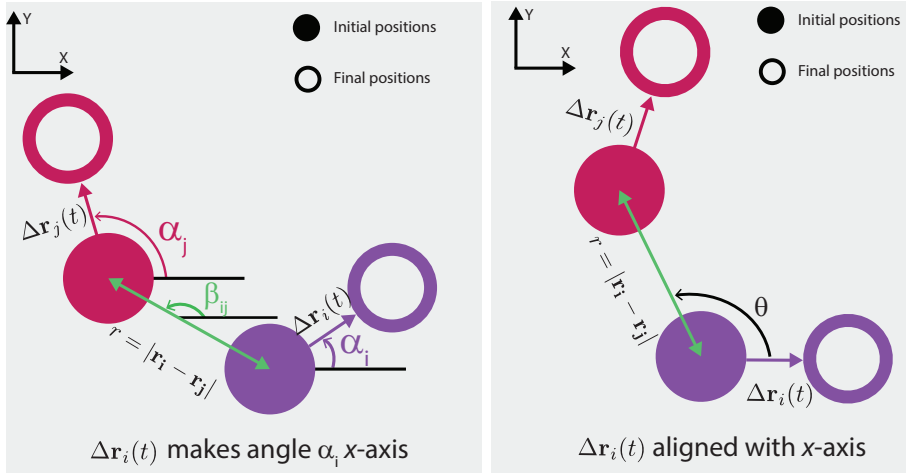


Fig. S7: Schematic for evaluating hydrodynamic displacement flow field. Consider a pair of particles, $\{i, j\}$. $\Delta \mathbf{r}_i(t = 0.5s)$, $\Delta \mathbf{r}_j(t = 0.5s)$, and the line joining i to j subtend angles α_i , α_j , and β_{ij} with respect to the positive x -axis (left image). To derive the plot shown in Figure 1f, we rotate the 2D coordinate system through an angle α_i for each reference particle i , so that $\Delta \mathbf{r}_i(t = 0.5s)$ is aligned along the positive x (horizontal) direction. The position coordinate of the j^{th} -particle is now defined as $\mathbf{r}(r, \theta)$ with i^{th} -particle as reference (right image).

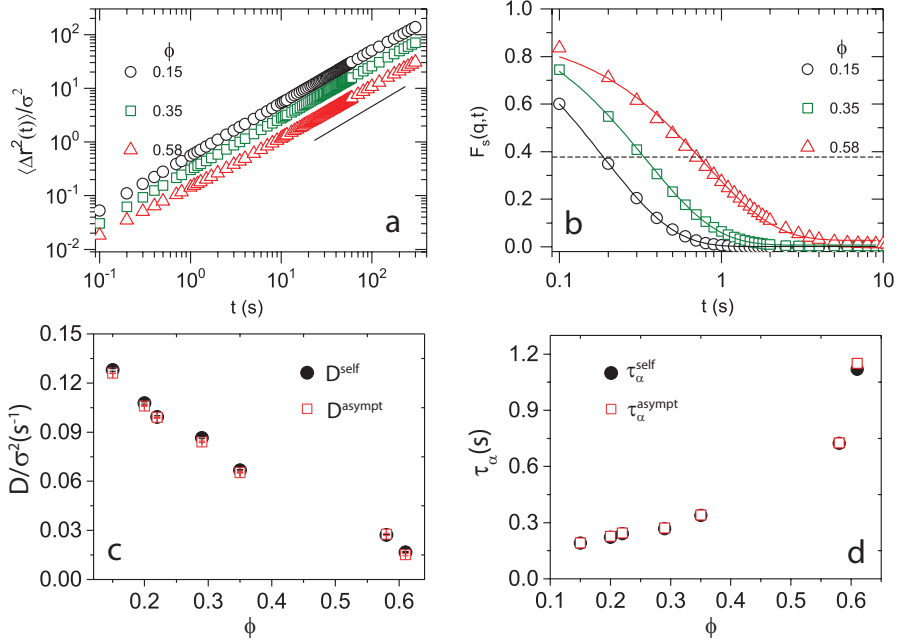


Fig. S8: Dynamics measurements. **a**, Mean squared displacement, $\langle \Delta r^2(t) \rangle$, and the **b**, self intermediate scattering function, $F_s(q, t)$, versus lag time t for particles in the samples at different packing area-fraction. Distinct colors (and shapes) represent different ϕ . D^{self} is measured from the linear regime in **a** shown by the solid line with slope 1. The solid lines in **b** are guide to the eye; relaxation time τ_α^{self} is defined from $F_s(q, t = \tau_\alpha^{self}) = 1/e$ (see horizontal dashed line). **c**, Comparison of D^{self} obtained from the data above (black symbols) with asymptomatic values derived in the same samples from the body frame analysis (red symbols), *i.e.*, $D^{asympt} = (D^L + D^T)$ for $r > 8\sigma$ as a function of ϕ . The errors bars in D are from fittings. **d**, Analogous comparison of τ_α^{self} (black symbols) with asymptomatic values obtained from the body frame (red symbols) $\tau_\alpha^{asympt} = (\tau_\alpha^L + \tau_\alpha^T)/2$ for $r > 8\sigma$ with ϕ .

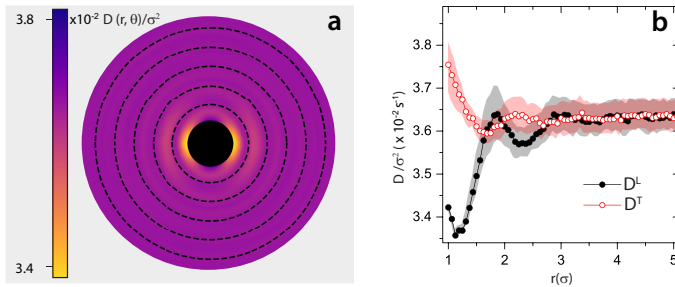


Fig. S9: Spatially inhomogeneous and anisotropic diffusivity at $\phi = 0.35$. **a** Polar colormaps $D(r, \theta)$ versus r . The dashed radial circles are at $r = \{2, 3, 4, \dots\}\sigma$. **b**, $D(r)$ along L and T directions corresponding to $\theta = 0^\circ$ and $\theta = 90^\circ$, respectively. The error bars in D are from fittings.